

Neitzke I:

Joint w/ D. Gaiotto, G. Moore

- A ^{new} scheme for constructing HK metrics on the total spaces of \mathbb{C} -int. systems (w/ extra data $\Omega(Y)$ that obeys a well-known formula)

- By twistor theory (Hitchin, HKLR)

HK mfold means

" \mathbb{C} sympl. mfold over $\mathbb{C}P^1$ " (of a special sort).

ξ

Restricted $\xi \in \mathbb{C}^*$,

our construction will look like gluing together patches which are open subsets of $(\mathbb{C}^*)^{2r}$

[cf. Kontsevich-Schubert, Gross-Siebert, Auroux, Gross-Hacking-Koelb, ...]

Novelty: keep track of ξ -dependence.

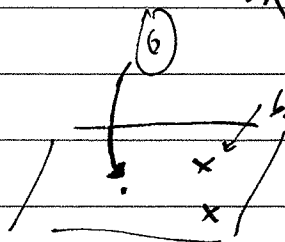
(fibers over 0, ∞ very different, not gotten from particularly the asymptotics as $\xi \rightarrow 0, \infty$)

Our HK metric will take the form

$$g = g^{sf} + (\text{instanton corr.})$$

simple, explicit.

weighed by Π 's of $\mathcal{R}(\delta)$



has fibers, $r=2$, g^{sf} does n't extend over bad fibers.

Need instanton corrections for that

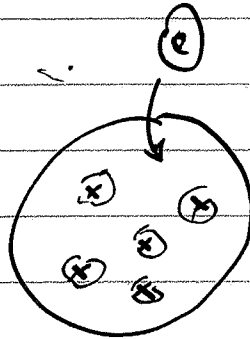
fix bad behavior of g^{sf} near sing. fibers.

Gross-Wilson:

elliptically fibered K3
generic \rightarrow 24 I_2 fibers

$$g = g^{sf}(R)$$

$R \rightarrow \infty$: torus fibers
shrink.



Now each bad fiber, G-W give us

"Ooguri-Vafa metric" corresponds to a single instanton correction,
exponentially close as $R \rightarrow \infty$

Plan: 1) Twistor geometry of $\mathbb{R}^2 \times T^2$

Ooguri-Vafa

2) general instanton corrections (KS, WCF)

3) Hitchin systems

HK geometry:

(M, g) Kähler w.r.t. \exists sympl. str. ω_i .
 \mathbb{C} str. J_i .

In fact, an S^2 worth of J 's: any $\sum a_i J_i$, $\sum a_i^2 = 1$.

Better, think of this as $\mathbb{C}P^1 \ni \mathcal{S}$

$\forall \mathcal{S}, M^{(\mathcal{S})} = (M, J^{(\mathcal{S})})$ is hol. sympl.

$$\omega := \frac{-i}{2\zeta} \omega_+ + \omega_3 - \frac{i}{2} \zeta \omega_-$$

$$[\omega_{\pm} = \omega_1 \pm i\omega_2]$$

enough to construct HK metric.

Take $\mathcal{M} = \mathbb{R}^2 \times T^2$, w/ flat product metric

determined by $\tau \in \mathbb{H}/SL(2, \mathbb{Z})$, $\text{vol}(T^2) = 1/R$

Fix \mathbb{C} -str. on \mathbb{R}^2 : \mathbb{C} coordinate a .

The HK str. we're after:

$$\mathcal{M}^{(s=0)} \simeq \mathbb{C} \times T^2$$

$$\mathcal{M}^{(s \in \mathbb{C}^*)} \simeq \mathbb{C}^* \times \mathbb{C}^*$$

Fix angular coords $\theta_e, \theta_m \in \mathbb{R}/2\pi\mathbb{Z} : T^2 \rightarrow \mathbb{R}/2\pi\mathbb{Z}$.

$$y = \theta_m - \tau \theta_e \quad y \sim y + 2\pi \quad \text{periodicity}$$

$$y \sim y + 2\pi\tau$$

then define

$$\chi_e = \exp[\pi R s^{-1} a + i\theta_e + \pi R s \bar{a}]$$

(complexify in a way dep. on s)

$$\chi_m = \exp[\pi R s^{-1} \tau a + i\theta_m + \pi R s \tau \bar{a}]$$

$$(\chi_e, \chi_m)(s): \mathcal{M} \rightarrow \mathbb{C}^* \times \mathbb{C}^*$$

$$[s \in \mathbb{C}^*]$$

[Pull back \mathbb{C} -str. $J^{(s)}$ from $\mathbb{C}^* \times \mathbb{C}^*$ to $M^{(s)}$]
 & hol. sympl. str. $d \log X_e \wedge d \log X_m$

$$\omega = \frac{1}{4\pi^2 R} d \log X_e \wedge d \log X_m$$

$$= \frac{-i}{2s} \omega_+ + \omega_3 - \frac{i}{2} s \omega_-$$

$$\omega_+ = \frac{1}{2\pi} da \wedge dy \quad (\text{hol. sympl. str. on } M^{(s=0)})$$

(need to rescale as $s \rightarrow 0$).

$$\text{i.e. } M^{(s \rightarrow 0)} \simeq \mathbb{C} \times T_{\mathbb{C}}^2.$$

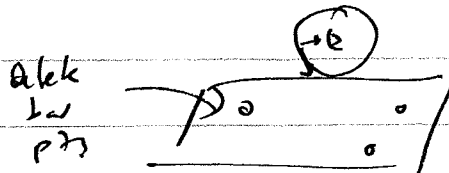
$$\omega_3 = \frac{i}{4\pi} \left[2\pi R (\text{Im } \tau) da \wedge d\bar{a} + \frac{1}{\pi R (\text{Im } \tau)} dy \wedge d\bar{y} \right]$$

$$X_e = \exp \left[\frac{1}{R} S \tau a + i \left(\theta_e + \pi R B a \right) \right]$$

now X_e, X_m have essential singularity as $s \rightarrow 0$.

idea: produce more interesting fns. w/ this essential singularity as $s \rightarrow 0, \infty$.

When have sol'n of J. (f' l) equations that X_e, X_m satisfy, have Stokes phenomenon.



Semiflat geometries:

Replace \mathbb{C} by a (1-dim '1) \mathbb{C} -unfold B'

Replace $(a, \tau a)$ in $X_{e,m}$ by general hol. fns. Z_e, Z_m :
 $B' \rightarrow \mathbb{C}$

$$\text{Now } \chi_e = \exp \left[\pi R z_e / \gamma + i\theta_e + \pi R \gamma \bar{z}_e \right]$$

Same construction \rightarrow simple HK metric

$$\text{At } M(\gamma=0), \text{ torus fibers have } \tau = \frac{dz_m}{dz_e}$$

for $\gamma \neq 0$, locally identified w/ $\mathbb{C}^* \times \mathbb{C}^*$.

(need assumption $\text{Im } \tau > 0$)

Ex: Take $B' = \{0 < |u| < 1\}$

$$z_e(u) = u$$

$$z_m(u) = \frac{1}{2\pi i} (u \log u - u).$$

(pick a cut, single-valued)

Same constr. as before:

\rightarrow HK metric g^{sf} on T^2 -bundle over B' .

but no way of extending over $u=0$

To extend HK str. over $u=0$:

Call our previous (χ_e, χ_m) χ_e^{sf}, χ_m^{sf} .

Improved version:

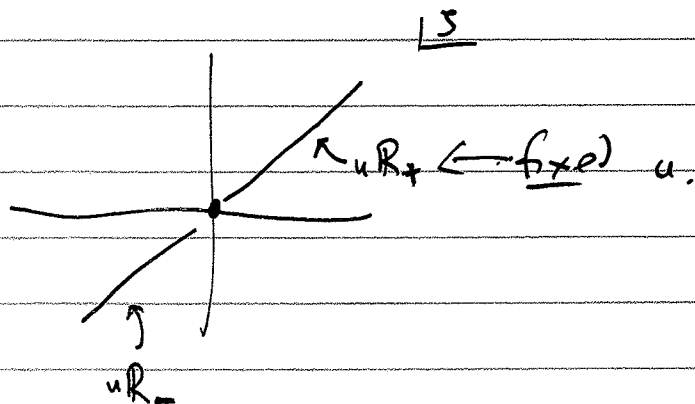
$$\chi_e = \chi_e^{sf}$$

$$\chi_m^{(\gamma)} = \chi_m^{sf(\gamma)} \exp \left[\frac{i}{4\pi} \int_{\mathbb{R}_+} \frac{d\gamma'}{\gamma'} \frac{\gamma'+\gamma}{\gamma'-\gamma} \log [1 - \chi_e(\gamma')] \right]$$

$$\chi_m^{(\gamma)} = \frac{i}{4\pi} \int_{\mathbb{R}_+} \frac{d\gamma'}{\gamma'} \frac{\gamma'+\gamma}{\gamma'-\gamma} \log [1 - \chi_e^{-1}(\gamma')] \Bigg]$$

$\chi_m(\zeta)$

$\chi_m(\zeta)$ is discontinuous as fun. of (u, ζ) :



As ζ crosses uR_+ , $\chi_m \rightarrow \chi_m(1 - \chi_e)$

" " uR_+ , $\chi_m \rightarrow \chi_m(1 - \chi_e^{-1})^{-1}$

[comes from $\Omega(\chi_e) = 1, \Omega(\chi_m) = 0$]

Despite this, $\omega = d \log \chi_e + d \log \chi_m$

is perfectly smooth, and

$$\text{still has } \omega = \frac{1}{z} \omega_+ + i \omega_3 + \zeta \omega_-$$

The correction to χ_m is finite in the limit $\zeta \rightarrow 0$.

("good asymptotics").

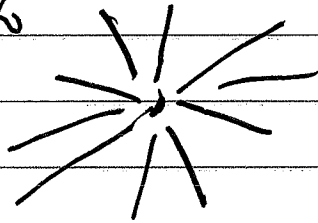
This χ_m $\xrightarrow{\zeta=0}$ simple asymptotic form $\sim \exp(\pi R \zeta_n / \zeta) \ll$,
exhibits Stokes phenomenon.

"analytic continuation of asymptotics \neq asymptotics of analytic continuation"

χ_m has Stokes phenomenon!

Can calculate metric explicitly: Doguri-Vafa nets:
only ingredient is the symplectomorphism.

Def: generalize to



pile of Stokes factors

when these rays hit each other, get W-C-F.