

# Day 4 Talk 4: Discussion

Possible topics:

- ① Twisted complexes and what does "derived" mean?
- ② Signs (of madness)
- ③ Grading on  $H\tilde{F}^*$
- ④ open-closed string theory
- ⑤ obstructions & curvature
- ⑥ Vasyl duality and mutations. / has to do w/ Lefschetz fibrations
- ⑦ immersed Lagrangians

Curved  $A_\infty$ -structures:

$Ob A$

$Hom_A(X, Y)$

$$\mu_A^d : Hom_A(X_{d-1}, X_d) \otimes \dots \otimes Hom_A(X_0, X_1)$$

Branchi  $\downarrow$

$$Hom_A(X_0, X_1)[2-d] \quad (d \geq 0)$$

$$\mu_A^1(\mu_A^0) = 0 \in Hom^3(X, X)$$

$$\mu_A^1(\mu_A^1(x)) = \mu_A^2(\mu_A^0, x) \pm \mu_A^2(x, \mu_A^0)$$

Ex:  $E \rightarrow M$  vector bundle w/ connection

$$A = \Omega^*(M, End(E))$$

$$\mu^1 = d \quad \mu^2 = \wedge, \quad \mu^0 = F_D$$

Ex 2:  $E \rightarrow M$  v.b. over a complex manifold

$$\mathcal{A} = \Omega^{0,*}(M, \text{End}(E))$$

$$\mathcal{A}^1 = \bar{\partial}\mathcal{A}, \quad \mathcal{A}^0 = F_{\bar{\partial}}^{0,2}$$

$d^2 = 0$  has interesting solutions (thanks to structure of nilpotent matrices)

$d^2 = 1$  less interesting solutions (get a splitting).

$(\mathcal{A}^0 = 0 \iff \text{vector field vanishes at } 0$   
(thinking about these as calculations)

But maybe the v.f. vanishes somewhere else!  
try to change/deform (infinitesimally) to extract something meaningful.

Technical condition:  $\mathcal{A}$  defined over  $\mathbb{Z}[\epsilon]$ ,  $\mathcal{A}^0$  of order  $\epsilon$ . (use formal deformation theory to make it go away).

$\mathcal{A}$  (obstructed)  $\Rightarrow$   $\tilde{\mathcal{A}}$  (unobstructed)  $\downarrow$  (similar to twisted complex)

$$\text{Ob}(\tilde{\mathcal{A}}) = \{(X, \alpha) \mid X \in \text{Ob } \mathcal{A}, \alpha \in \text{Hom}_{\mathbb{C}}^2(X, X),$$
$$\mathcal{A}^0 + \mathcal{A}^1(\alpha) + \mathcal{A}^2(\alpha, \alpha) + \dots = 0$$

inhomogeneous Maurer-Cartan  $\quad \quad \quad = \mathcal{A}^2$

Alternatively, just consider category of all such objects, & throw away obstructed guys.

(FOOD call this "filtered" ASD alg.)

You can do this & twisted complexes at once, but not related.

Version of this you can do which hasn't really been written down.

If connection is prolocally flat, induced connection

on  $T^*M$  bundle is flat -

(think of  $H^1$  as being central).

$$p \in C^2(A, A)$$

$$dp = 0 \text{ of order } k$$

consider  $\tilde{A}_{[p]}$  (not curved)

$$Ob(\tilde{A}_{[p]}) = \left\{ (x, \alpha) \mid \left[ \eta^0 + \eta^1(\alpha) + \eta^2(\alpha, \alpha) + \dots \right] \right\}$$

Why is this important?

$$= \rho^0$$

$$(M, \omega) \quad c_1(M) = 0$$

small, b/c no hol. disks of area  $> 0$   
↳ (power controlled by  $\omega(B)$ ).



$\eta^0 =$  boundaries of hole-discs

Try to solve

$$\eta^0 + \eta^1(\vec{a}) = \dots$$

gradually, step by step.

$$\underline{\text{Ex:}} \quad L^3 \subset M^6, \quad c_1(M) = 0$$

$$H^1(L) = H^2(L) = 0$$

$\Rightarrow \exists$  unique object in  $\mathcal{F}(\dots)$

(i.e. can solve, & solve uniquely)

For other cases, generally an infinite-dim'l space of paths/choices

physics: Objects  $(L, A \in \Omega^1(L, \mathbb{R}))$ ,

$$\omega|_L = 0, \quad F_A = 0 \quad (\text{to guarantee good conditions in twisted or male})$$

this is before instanton corrections.

$$\text{Actually, what's true is } \frac{i}{2\pi} F_A = -PD(\eta^0).$$


Max's examples

fun w/  
singularities. ↑  
current

$\rho = \text{const.} \cdot \text{id}$ , so can translate away.

( $c_1(M) \neq 0$ , so not graded any more)


⑦ Immersed Lagrangians (Joyce, -)

  $L \xrightarrow{\tau} M$  (Lagrangian Immersion)  
 (lots of examples & existence  
 thms).

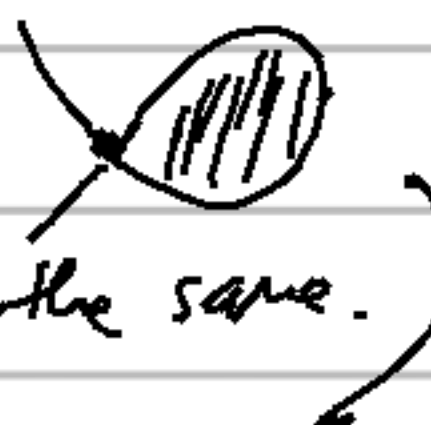
$$CF^*(L, L) = C^*(L) \oplus \bigoplus_{(xy) \in L} \mathbb{Z}\langle x, y \rangle$$

$\uparrow$   
 Morse complex       $\tau(x) = \tau(y)$   
 $x \neq y$

(why ordered pairs?  $\nearrow$ )


 when you flow, 2 intersection points!

$\eta^0 =$  holo discs +



claim: everything else is the same.

1. e.  $u: \mathbb{H} \xrightarrow{J\text{-hol.}} M$        $\lim_{s \rightarrow +\infty} u(s) = x$   
 $\tilde{u}: \mathbb{R} \rightarrow L$        $\lim_{s \rightarrow -\infty} u(s) = y$   
 $\tau(\tilde{u}) = u|_{\mathbb{R}}$

ex:  seems unobstructed to naked eye, but look at differential.

Usually, think of all solutions to Maurer-Cartan as arising from isodropics?

## Signs and gradings

$$L_0, L_1 \subset \mathfrak{M}$$

$$\mathcal{P} = \{x: [0, 1] \rightarrow \mathfrak{M} \mid u(0) \in L_0, u(1) \in L_1\}$$

$$\mathcal{P} \rightarrow \mathcal{U}_\infty / \mathcal{O}_\infty$$

①  $H^1(\mathcal{U}_\infty / \mathcal{O}_\infty) = \mathbb{Z} \rightarrow H^1(\mathcal{P})$  obstruction

②  $H^2(\mathcal{U}_\infty / \mathcal{O}_\infty; \mathbb{Z}/2) = \mathbb{Z}/2 \rightarrow H^2(\mathcal{P}; \mathbb{Z}/2)$  obstruction  
to grading  
to signs.

①  $2c_1^{\text{rel}} \in H^2(\mathfrak{M} \times [0, 1], L_0 \times \{0\} \cup L_1 \times \{1\})$

↓ transgression

$$H^1(\mathcal{P})$$

↑ try to kill this

②  $w_2(L_0) \oplus w_2(L_1) \in H^2(L_0) \oplus H^2(L_1)$

↓ evaluation

$$H^2(\mathcal{P}; \mathbb{Z}/2)$$

(Some version of fully Atiyah-Singer index theorem.)

① some sort of Calabi-Yau condition, but more.

Killing 2c, :

$$(M, \omega, J) \quad (\Lambda_{\mathbb{C}}^n TM)^{\otimes 2} \stackrel{\text{fix}}{=} \underline{\mathbb{C}}$$

("you don't want to know it's dead, you want to just kill it explicitly")

Let's say we're killing  $c_1$ : choose a complex n-form  $\eta \in (\Lambda_{\mathbb{C}}^n TM)$ .

$$L \text{ oriented } \subset M \rightarrow \eta_L \in H^1(L)$$

$$\eta_L = [\alpha_L : L \rightarrow S^1], \quad \alpha_L(x) = \eta(TL_x)_{\text{phase}}$$

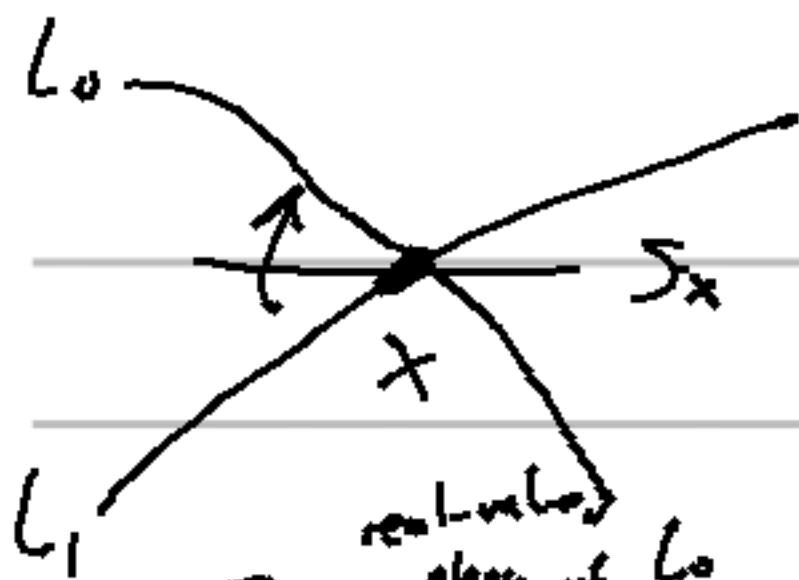
Choose a real-valued phase, a lift

$$\tilde{\alpha} : L \rightarrow \mathbb{R} \text{ of } \alpha.$$

M surface: trivialization of  $TM =$  oriented foliation,  $F$



$\alpha_L$  gives you a way of rotating  $TL$  into the foliation.



$$F_x \xrightarrow{\text{real-valued phase of } L_0} (TL_0)_x \xrightarrow{\text{"short" map}} (TL_1)_x \xrightarrow{\text{real-valued phase of } L_1}$$

count winding #

Thomas-Yau: apply this to uniqueness of special Lagrangian.

This reproduces: the classical grading in Morse theory.

signs: Pick a spin structure (not an ant class, an actual spin structure)  
 spin automorphism  $(-1)$  acts on HF by  $(-1)$  — very messy.

Families index theory

$$\mathcal{P} \rightarrow U/O$$

$$\Omega \mathcal{P} \rightarrow \Omega(U/O) \cong \mathbb{Z} \times BO_{\infty} = \text{Fredholm operators}$$





loop in  $\mathcal{P}$ .

$\bar{D}$ -op. on  $u^*TM$   
with boundary values  
on  $L_0, L_1$ .

$$u: [0, 1] \times S^1 \rightarrow M.$$

compute using Atiyah-Singer

Next obstruction:

$$H^2(\mathcal{P}; \mathbb{Z}/2) \rightarrow H^1(\Omega\mathcal{P}; \mathbb{Z}/2) \leftarrow H^1(\mathbb{Z} \times B\mathbb{O}_\infty, \mathbb{Z}/2)$$

determinant line of Fredholm  
operator

$\cong \mathbb{Z}/2$ .

$w_1$  (det line bundle)

If  $u$  varies, do you have a 1-param. family  
of det. lines, what is the det. line bundle?

trivial or not? family index thing.

(Atiyah-Singer pt. 5)

# Open-Closed String Theory

How does  $F(-)$  fit into an O-C str. theory?

## Feynman theory:

Open-closed: framed little disc operad.

$$\text{Diagram: } \mathcal{O}(1) = \left\{ \text{Diagram: } \bigcirc \right\} \cong \mathcal{S}^1$$

$$\mathcal{O}(2) = \left\{ \text{Diagram: } \bigcirc \right\} \cong T^3$$

$$\mathcal{O}(3) = \dots$$

depend by stacking disc  $\mathcal{O}$  putting it, aligning marked points.

Algebra over  $H_{\text{str}}(\mathcal{O})$ : BV-algebra

$V$  graded vector space, differential


$$\Delta: V \rightarrow V[-1]$$

$V \times V \rightarrow V$  commutative, associative.

$\mathbb{Z}^3$  classes give  $x \cdot y, x \cdot \Delta y, \Delta x \cdot y, \Delta(x \cdot y);$   
 $(\Delta^2 = 0).$

The failure of  $\Delta$  to be a derivation, is a Lie bracket;

hp.  $\{x, y\} = \Delta(x \cdot y) - x \cdot \Delta y - \Delta x \cdot y$

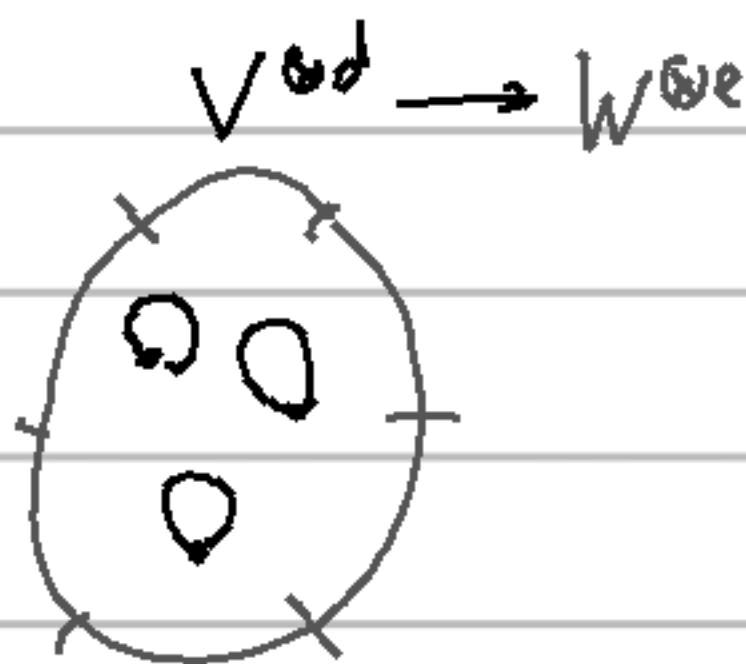
Jacobi from  $\mathcal{O}(S) =$  

This is a genus 0

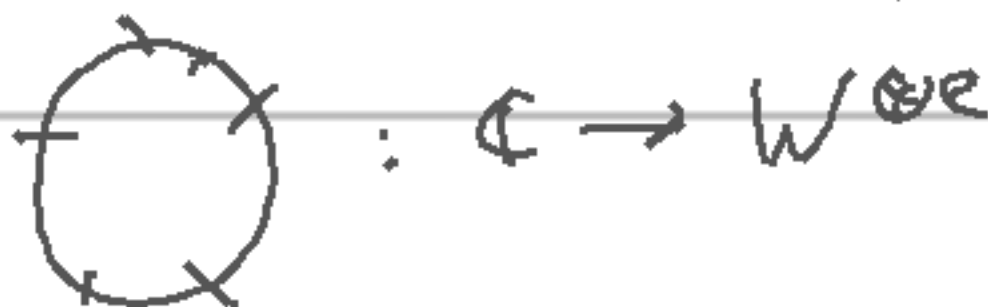
closed string topological field theory (i.e.

a BV algebra)

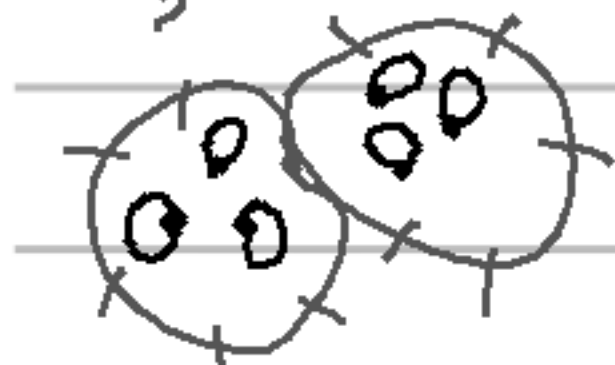
open-closed



$V, W$  carries a graded symmetric pairing



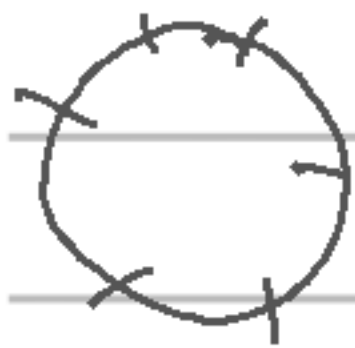
compactification



open-closed top. string theory.

\* bricolage,  $(\Sigma, \text{cyclic})$  operad

↑ puncture boundary pts.



$: \mathbb{C} \rightarrow W \otimes \mathcal{L}$

gives  $W$  the structure of an  $A_{\infty}$  algebra which is cyclic  $(\mathbb{C} - \gamma)$

(Fuk categories also have this, up to homotopy, after lots of work)



homotopy BV-algebra

can look at chains of these guys, but get no new info. (some sort of formality result)

mixed guys:



what happens to these?



gives a map of homology BV algebras (similar to Sheel's talk).

$$(x) \quad \mathcal{V} \rightarrow CC^*(W, W)$$

$\swarrow$   $\uparrow$  (it has structure of BV algebra, requiring some work)

$$HH^2(W, W)$$

carries  $\Delta$  = Carnes' operator, Sheel's product.

should be a formal consequence of product structure of these moduli spaces.

$$\left( \begin{array}{l} \text{idea: } C^*(M) \\ \parallel \\ \mathcal{V}^{\text{odd}} \rightarrow W^{\text{odd}} = CF^*(L, L) \end{array} \right).$$

Can recover some closed stuff, stuff from  $W$ ,

i.e. set  $\mathcal{V} = CC^*(W, W)$   $(x) = \text{id}$ .

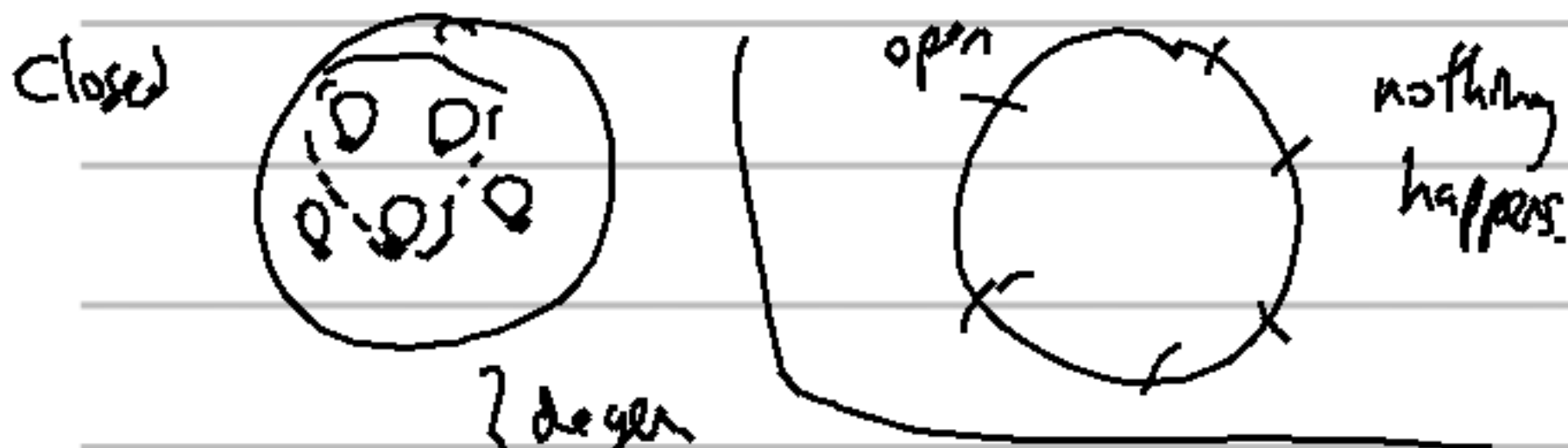
(don't require  $(x)$  to be a quasi-iso, i.e.  $W=0$ )

but there are classes in which this is the case

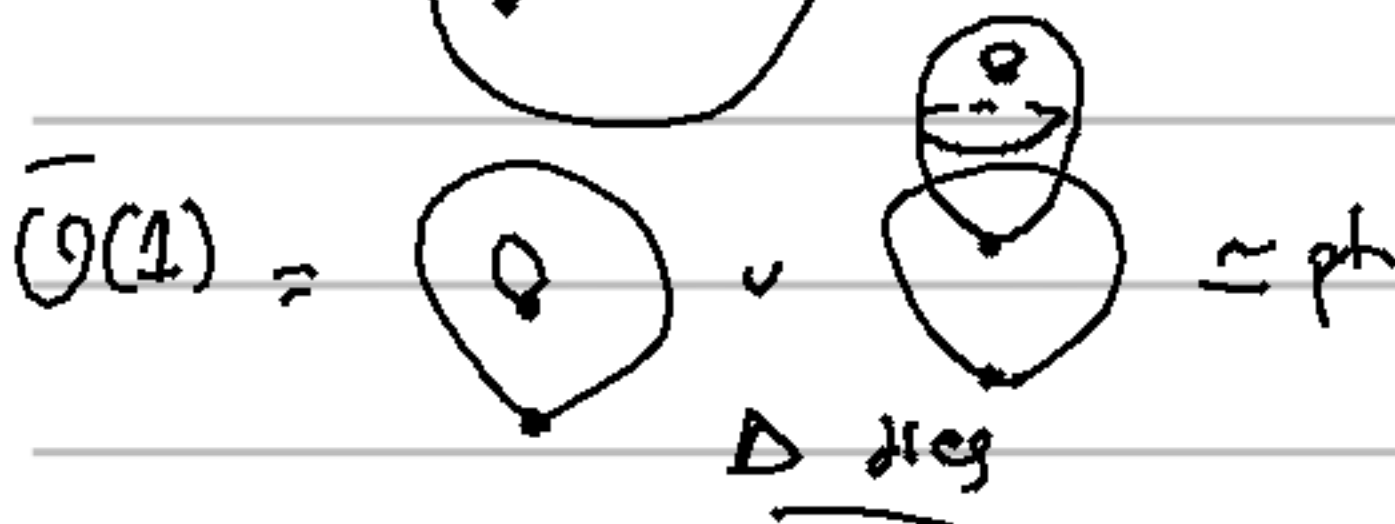
(open-closed TFT := alg. over chains on bicovariant algebra.)


# What actually happens (impress)


Moduli spaces are further compactified to Deligne-Mumford spaces



closed-string  
 Open  $\bar{\mathcal{O}}$ .



$\bar{\mathcal{O}}(2) =$    $\cong \text{pt}$  - (consider all degenerations)  
 need to

$\bar{\mathcal{O}}(3) \cong \bar{M}_{0,4} \cong S^2$  (Paul's claim). 

product

genuinely new 3-fold product

open-string: cyclic  $A_\infty$ -structure.  $V$

closed-string sector: Coh FT (chronological field theory)  $W$

Kontsevich - (so of generators) in general, some relations between product, Manin, Dubrovin, Givental) not enough.



open-closed  
string sector



still have  $V \rightarrow C^*(W, W)$   
Coh FT Homotopy BV

If this is a quasi-isom.  $\Rightarrow$  Connes boundary operator vanishes

$\Rightarrow$  spectral sequence from  $H^*(W, W)[u]$  to  $H^*(C^*(W, W))$  degenerates (Connes  $bd.$  is first differential).

We would like the spectral sequence  
to degenerate anyway i.e., get a  
map of Hodge CohFT's.

Kontsevich conjecture: If  $A_\infty$  structure "smooth,"  
get this degenerates.

Arithm from Kähler geometry; degeneration of  
Hodge-de Rham forms...

On  $\bar{F}$ , we have all these structures, in particular.

$$\Delta \rightarrow 0 \Rightarrow \{ -, - \} = 0 \text{ on } V.$$

i.o.  $V \rightarrow \mathbb{C}^*(W, W)$  map of dgla's  
back to 0.

$\Rightarrow$  family of  $A_\infty$ -algebras over  $V$   
(big Fukaya category).

(really useful for, e.g. toric varieties).

i.e. get a sheaf of  $A_\infty$  categories over  $V$ , fiber  
at 0 is usual Fukaya category.



$C_n$ , e.g., deform by  $C_1$ , which has some geometric significance.

As a standard deformation

deg  $k$  piece  $\rightarrow$  rescale by  $\lambda^k$ .  
suppresses higher order terms

on the whole, unsatisfactory.

Can't really reconstruct  $V$  from  $C^*(W, W)$

In principle, know what you have to do to upgrade to topy BV  $\rightarrow$  Coh FT.

Additional piece of info required (to do w/ killing  $\Delta$ ). First manifestation is a power series in two variables w/ values in  $HC^*(W, W)$ .

Given  $W$ , have different choices of data, give different Coh FT's. Group that acts on those is  $\text{the } \infty\text{-dim'd, like a loop group of a symplectic } \mathbb{R}$   
"Gervais's twisted loop group." cf. Kevin Costello.

Costello knows how to add in the extra data,

there is absolutely a higher gen-5 version of this.